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Forward-stepwise regression analysis for fine leak batch testing of wafer-level hermetic MEMS packages

Changsoo Jang^a, Byeng Dong Youn^a, Ping F. Wang^a, Bongtae Han^{a,*}, Suk-Jin Ham^b

^a University of Maryland, College Park, MD 20742, USA

^b Samsung Electro-Mechanics, Central R&D Institute, Suwon, Gyeonggi-do 443-743, Korea

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ABSTRACT

An advanced regression scheme is proposed to analyze fine leak batch testing data of multiple MEMS packages. The scheme employs the forward-stepwise regression method to infer the information of leaky packages from a batch test data. The analysis predicts the number of leaky packages and the true leak rate of each leaky package in a progressive manner. The scheme is implemented successfully using an actual batch test data obtained from wafer-level hermetic MEMS packages. An error analysis is followed to define the applicable domain of the scheme. Advanced formulations are also suggested to extend the applicable domain.

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1. Introduction

The helium mass spectrometer based leak test is widely used in the industry for fine leak detection [1,2]. The output of the spectrometer is an apparent leak rate, which is defined as the leak rate of a given package as measured under specified test conditions [3]. In practice only the initial apparent leak rate is used as a *qualitative* measure of hermeticity.

An individual true leak test is essential to inspection of the hermetic performance of each package. The information of individual leaky packages is often required when a fabrication process is to be evaluated. However, such an individual test is not practical in a mass production environment because of long testing time. A test with multiple packages is a preferred choice in the industry, where a batch of packages is subjected to the same test conditions. Yet, the current procedure of batch tests can provide only qualitative information about possible leaky packages in the batch, which is not sufficient to address the issues associated with manufacturing yields.

In this paper we extend the He fine leak test based true leak rate measurement into the batch test domain. The scheme employs the forward-stepwise regression method to determine the number of leaky packages and their true leak rates from a single batch test data.

2. Background: gas leak rate of an individual package

In the metallic seals of wafer-level hermetic MEMS packages, gas transport occurs through randomly present nano-scale leak

* Corresponding author. Tel.: +1 3014055255.

E-mail address: bthan@umd.edu (B. Han).

channels and thus is governed by the molecular gas conduction theory [4]. Even when multiple nano-scale leak paths exist, they can be modeled as an effective single leak channel. The leak rate depends on the gas molar mass and the geometry of the channel (diameter and length). The conduction-based governing equation for leakage, known as Howl–Mann equation [5], is expressed as

$$R_{i} = \frac{LP_{b}}{P_{0}} \left[1 - \exp\left(-\frac{Lt_{b}}{VP_{0}}\right) \right] \exp\left(-\frac{Lt_{dwell}}{VP_{0}}\right)$$
(1)

where R_i is the *initial* apparent leak rate obtained at the beginning of spectrometer operation (atm-cc/s); L is the true leak rate of a gas (atm-cc/s); t_b is the bombing time (s) during which a package is subjected to helium pressurized at the bombing pressure, P_b ; t_{dwell} is the dwell time (s) which is the short duration between the instant the specimen is taken out of the bombing chamber and the instant the spectrometer is switched on; V is the package cavity volume (cc); and P_0 is the standard pressure (=1 atm). The existing standard [3] utilizes the equation for hermeticity qualification. It should be noted that the square root ratio between molar masses of helium and air, which is included in the original Howl–Mann equation to convert a helium leak rate to a true air leak rate (L_a), is omitted in Eq. (1) because only the helium leak rate is considered in this study.

As can be seen from Eq. (1), the initial apparent leak rate can vary significantly with testing parameters. Therefore, the standard using the initial apparent leak rate can serve only as a qualitative benchmark for the small cavity volumes.

Recently, the fine leak test has been advanced to extract the true leak rate of a single micro to nanoliter package quantitatively from the complete apparent leak rate history, i.e., R(t) for $t > t_{dwell}$ [4,6]. Two new stages after the dwell time [6] were introduced in the procedure: (i) the zero signal time during which a small





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amount of helium present in ambient air inside the test chamber of the spectrometer is drawn out, and (ii) the measurement phase during which apparent leak rates are collected for a subsequent analysis. All four stage of the fine leak test is shown in Fig. 1 [6]. Using the original Howl–Mann equation, the apparent leak rate, Ω (atm-cc/s), at the end of the zero signal time, t_{zero} , can be expressed as

$$\Omega = \frac{LP_b}{P_0} \left[1 - \exp\left(-\frac{Lt_b}{VP_0}\right) \right] \exp\left[-\frac{L(t_{dwell} + t_{zero})}{VP_0}\right]$$
$$= R_i \exp\left(-\frac{Lt_{zero}}{VP_0}\right)$$
(2)

Then the apparent leak rate during the measurement phase can be simplified as

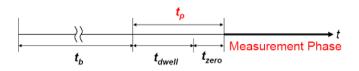
$$R(t) = \Omega \exp\left(-\frac{Lt}{VP_0}\right) \tag{3}$$

where *R* is the apparent leak rate (atm-cc/s) and *t* is the elapsed time during the measurement phase (s). The above equation describes well the apparent leak rate history since the gas flow inside a spectrometer chamber is always in the molecular conduction regime [4]. Then the true leak rate can be calculated accurately through a regression analysis [6].

Another important practical aspect of the fine leak test is the measurement range associated with the spectrometer sensitivity. Due to the finite helium spectrometer sensitivity, the fine leak test cannot detect true leak rates below or above certain limits. Eq. (2) is used to illustrate the limits. In Fig. 2, the initial leak rate is plotted as a function of true leak rates. The parameters used in the calculation were: $P_b = 5$ atm, $t_b = 6$ h, $V = 2.156 \times 10^{-4}$ cc, $t_{dwell} = 10$ min and $t_{zero} = 150$ s. The plot shows that the detectable true leak rate ranges from 4.5×10^{-10} atm-cc/s to 3.5×10^{-6} atm-cc/s if the spectrometer sensitivity is 10^{-10} atm-cc/s. In practice this range is reduced further due to measurement uncertainties associated with the data acquisition system.

3. Methodology

The idea of inferring the information about leaky packages from a batch test result is based on the fact that each leaky package generates a unique signal pattern defined by Eq. (3), and accordingly, the total signal should also have a unique pattern. A total apparent leak rate of *n* leaky packages (R_b) can be expressed as



- Bombing time, t_b: The duration for which the specimen is exposed to pressurized helium (on the order of hours)
- Dwell time, t_{dwell}. Time elapsed between the end of bombing and the start of the spectrometer (on the order of about 10 minutes).
- Zero signal time, *t_{zero}*: Time taken by the zero signal to stabilize (on the order of a couple of minutes).
- Preprocessing time, t_p: Time elapsed between the end of bombing and the point at which the zero signal stabilizes;
 i.e., t_p = t_{dwell} + t_{zero}

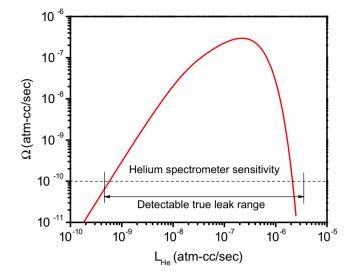


Fig. 2. Initial apparent leak rates as a function of true leak rates (cavity volume = 2.156×10^{-4} cc, bombing time = 6 h, dwell time = 10 min, zero signal time = 150 s).

$$R_b(t) = \sum_{i=1}^n \Omega_i \exp\left(-\frac{L_i t}{V P_0}\right) \tag{4}$$

The unknowns in the above equation are the number of leaky packages (n), and the initial leak rate (Ω_i) and true leak rate (L_i) of each leaky package. This total signal can be decomposed into the original set of individual signals using a proper data processing technique. The forward-stepwise regression method is employed to conduct this multi-parameter, non-linear inverse analysis.

A forward-stepwise regression analysis [7] uses an additive form of exponential functions [Eq. (4)] to model the apparent leak rate data. The analysis predicts the number of leaky packages (n), the initial leak rate $(\Omega_i, i = 1 \text{ to } n)$, and the true leak rate $(L_i, i = 1 \text{ to } n)$ in a progressive manner. The regression analysis begins with n = 1 and increases n gradually until a convergence criterion is met.

3.1. Regression criterion

The prediction square sum error is used as a convergence criterion to ensure that the prediction accuracy is achieved to a desired level. The normalized L^2 regression error norm (e_I) is defined mathematically as

$$e_{I} = \sum_{t=0}^{t_{e}} \left(\frac{R(t) - \widehat{R}(t)}{\mu_{R}} \right)^{2} \leqslant \varepsilon_{I} \quad \text{and} \quad \mu_{R} = \frac{1}{n} \sum_{t=0}^{t_{e}} R(t)$$
(5)

where R(t) is the regression model of the apparent leak rate, n is the number of data points ($= t_e/\Delta t + 1$; t_e and Δt are the time duration and interval of data acquisition, respectively) and μ_R is the mean of the observed apparent leak rate.

In general, a smaller regression error norm would provide higher convergence accuracy. If it is too small, however, the analysis may not converge. The error norm should be determined carefully considering how well the mathematical function describes the physical phenomenon (i.e., the assumptions used in the mathematical model) and the level of noise in experimental measurements. More importantly, the non-linear regression analysis with multiple roots can produce values which are not physically valid. Extra criteria are required to reject invalid solutions.

3.2. Extra criteria for physical

The following extra criteria are proposed to avoid mathematically feasible but physically invalid roots and thus ensure the convergence of the regression analysis in a strong sense. The extra criteria are developed from three physical conditions that should not be violated; (1) the true and initial leak rates are positive, (2) the resolution of the helium spectrometer is limited and the apparent leak rate should be larger than the spectrometer sensitivity limit, and (3) there should be one-to-one correspondence between the true leak rate and the apparent leak rate. The first criterion is simple to implement and the criteria corresponding to the last two conditions are described below.

(1) Minimum value for the initial leak rate (e_{II}) :

$$e_{II} = \min(\Omega_i) > \varepsilon_{II} \tag{6}$$

This regression analysis takes the minimum Ω value as a convergence indicator for the second physical condition. When $\Omega_i \leq \varepsilon_{ll}$, a package is regarded as hermetic. The value of ε_{ll} should be the minimum initial leak rate that can be captured by the helium spectrometer.

(2) Minimum difference of the true leak rate values (e_{III}) :

$$e_{III} = \min[|L_i(t) - L_i(t)|] > \varepsilon_{III}, \quad \forall i, j \in [0, n]$$

$$(7)$$

The difference between any of two predicted values for a true leak rate can be used as another convergence indication for the third physical condition. Numerical noise in the apparent leak rate may lead to an artificial increase in the number of leaky packages (*n*) by numerically splitting one leaky package into multiple artificial ones with the same or numerically equivalent leak rate but different Ω values. The criterion $\varepsilon_{\rm III}$ should correspond to the minimum measurable true leak rate.

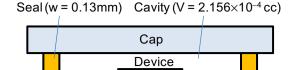


Fig. 4. Schematic of test MEMS package.

Substrate

A flow chart of the proposed regression analysis is shown in Fig. 3. The regression process ends when the regression criterion [Eq. (5)] is first met. If the second and third criteria are also met, the analysis is completed. If either of the extra criteria is violated, the initial values of the regression are changed and a new regression analysis is performed with the known number of package. This trial step is repeated until all the criteria are met or the number of trials reaches a certain number (j_e in Fig. 3). If the results do not satisfy all the criteria at the end of j_e trials, the analysis ends and the results at the previous step (n - 1 packages) becomes the final solution.

4. Implementation

The proposed method was implemented for an actual batch test data. The package and the experimental procedure are described briefly. The batch test results are presented first and the validation test results are followed.

4.1. Package and test procedure

The wafer-level hermetic MEMS package used in the test is shown schematically in Fig. 4. The cavity volume was 2.156×10^{-4} cc. The seal was made by the Sn–Au eutectic bonding process and the width of the metallic seal was 0.13 mm.

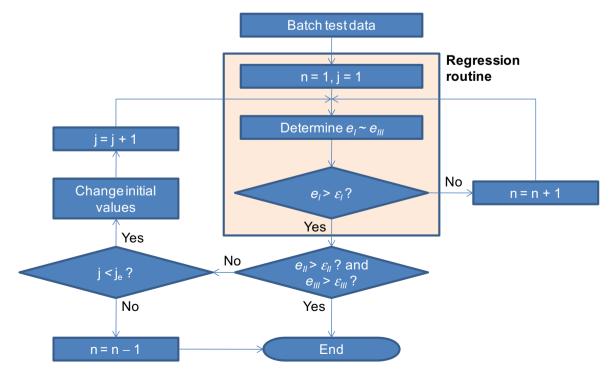


Fig. 3. Flow chart of the proposed regression analysis procedure.

We followed a standard leak test procedure described in Fig. 1, which was implemented for an individual package in Ref. [4]. In each leak test, the signals (apparent leak rates) were documented continuously for 1–2 h after 6 h of bombing at 5 atm and 10 min of dwell time for transferring specimens from the bombing chamber to the spectrometer chamber. Details about the fine leak test procedure and experimental setup can be found in [8].

Fig. 5 illustrates the entire test and regression procedure that was performed in this study. A total of 100 packages were fabricated and it was confirmed through the first-level batch test (Batch 1) that there should be leak package(s). The signal obtained from Batch 1 was used for a regression analysis and the information about leaky packages was attained. In order to verify the regression result, extra batch tests as well as individual tests were performed. In the second-level test the original batch was split into two subbatches (Batch 1-1 and 1-2). The sub-batches were subjected to the same leak testing and only those with evident leakage were selected for the next level batch tests. Such hierarchical testing continued until we obtained complete leakage information of each leaky package. As a result, 28 leak tests were carried out in total.

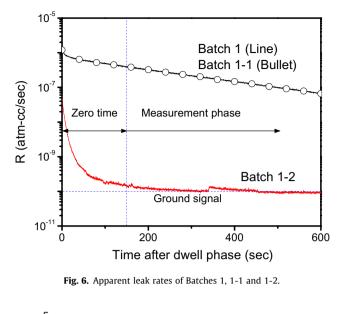
4.2. Result and analysis

The apparent leak rates of the first level and second level batches, documented immediately after the dwell time, are shown in Fig. 6 in the semi-log scale. Batch 1 (and Batch 1-1) contained leaky packages while no leaky packages were included in Batch 1-2. Due to the same content of leak packages, Batch 1 and 1-1 produced practically the same leak rate curves, confirming the repeatability of the helium leak test conducted in this study. It also confirmed that the contribution of hermetic packages to leak rate signal is negligible and thus Eq. (4) is valid.

The initial apparent leak rate is a sum of the zero signal (preexisting helium gas in the spectrometer chamber) and gas flowing out of leaky packages. The signal of Batch 1-2 was produced only by the zero signal and it was used to determine the zero signal time; it was determined to be 150 s after which the signal was stabilized to the ground level $(10^{-10} \text{ atm-cc/s})$ [8].

The signal of Batch 1 of the measurement phase (after the zero signal time) is plotted in Fig. 7 in the linear scale. It was used for the subsequent regression analysis. The three criteria required for the forward-stepwise regression were first determined for the package and spectrometer used in the test. An optimal value of ε_{l} was obtained as 1×10^{-8} , which was comparable to the absolute regression error $(10^{-4})^2$ or 0.01%. The values of ε_{ll} and ε_{lll} were determined as 10^{-10} atm-cc/s and 5×10^{-10} atm-cc/s, which correspond to the spectrometer sensitivity and the minimum measurable true leak rate, respectively.

The results of the regression analysis are presented in Table 1. Two leaky packages were identified in the batch. The true leak rates of the two packages were 5.90×10^{-6} and 8.50×10^{-7} atm-cc/s and the corresponding initial leak rates were 1.1×10^{-8} and 3.92×10^{-7} atm-cc/s. The regression results are compared with the experimental data in Fig. 6; they are virtually identical, indicat-



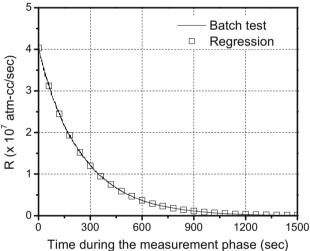


Fig. 7. Apparent leak rate signals of Batch 1 over the measurement phase (after the zero signal time) are plotted along with the results from the proposed regression analysis (in the linear scale).

ing that the regression was effective. It is to be noted that discrete regression data points are shown in Fig. 7 in order to distinguish them from the experimental data.

4.3. Validation

As shown in Fig. 5, Batch 1-1 was broken into five sub-batches, among which Batch 1-1-2 and 1-1-5 were identified to have leaky packages. Individual tests were conducted for all packages within the two batches. Packages #19 and #41 were finally found to be

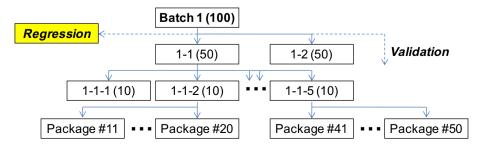


Fig. 5. Procedure of batch and individual leak tests (numbers in parentheses are the number of packages contained in batches).

 Table 1

 Regression results with batch and individual test data.

	Package 1	Package 2
Batch $L_i\left(\Omega_i ight) imes 10^{-7}$	59.0 (0.11)	8.50 (3.92)
Package #19 $L_i\left(\Omega_i ight) imes 10^{-7}$	49.5 (0.11)	
Package #41 $L_i\left(\Omega_i ight) imes 10^{-7}$		8.51 (3.93)
Difference (%)	19.2 (0.0)	0.1 (0.3)

leaky. Fig. 8 shows signals from individual and the corresponding sub-batches. They are overlapped well within the measurement sensitivity $(10^{-10} \text{ atm-cc/s})$, confirming again Eq. (2). The true and initial leak rates were determined through regression analyses. The true leak rates were 4.95×10^{-6} and 8.51×10^{-7} atm-cc/s and the corresponding initial leak rates were 1.1×10^{-8} and 3.93×10^{-7} atm-cc/s, for packages #19 and #41, respectively. The results of the regression are also plotted in Fig. 8.

The individual testing results are compared with the batch test results in Table 1. The batch test predicted the number of leaky packages accurately. While the initial leak rates of the two leaky packages were identical, one of true leak rates differed by 20% between batch and individual analyses. This discrepancy resulted from a considerably small apparent leak rate of package #19 compared to the measurement sensitivity and intrinsic experimental uncertainties. Although acceptable in practice, the cause of the discrepancy is analyzed through an uncertainty analysis in the following section, which defines the applicable domain of the proposed scheme.

5. Domain of application

As discussed earlier, the fine leak test has the limited application domain due to the finite spectrometer sensitivity. The same limitation should apply to the batch testing. In addition, the batch testing data should be handled more carefully due to potential multiple solutions (non-uniqueness) even when all leaky packages in a batch lie in the detectable true leak range. The following analysis is provided to define the extra constraints of the batch testing for solution uniqueness.

5.1. Error analysis

Let us consider a case where the combined signal of two leaky packages (packages 1 and 2) is similar to a signal from a single leaky package. It can be expressed as

$$\Omega \exp(-\lambda t) \approx \Omega_1 \exp(-\lambda_1 t) + \Omega_2 \exp(-\lambda_2 t) \quad \text{where } \lambda = L/VP_0$$
(8)

An "approximate-equal" sign is used in Eq. (8) to take account of the allowable error norm in the regression analysis. The Taylor series of the above equation can be written as

$$\Omega\left[1 - \lambda t + \frac{\lambda^2 t^2}{2} + \cdots\right] \approx \Omega_1 \left[1 - \lambda_1 t + \frac{\lambda_1^2 t^2}{2} + \cdots\right] + \Omega_2 \left[1 - \lambda_2 t + \frac{\lambda_2^2 t^2}{2} + \cdots\right]$$
(9)

In the proposed regression analysis, data obtained at the early stage of measurement phase are most important since the error norm deals with absolute (not logarithmic) difference between the data and prediction. Considering that $\lambda_i t$ at the early stage is typically on the order of $10^{-2}-10^{-3}$, it is reasonable to assume that the con-

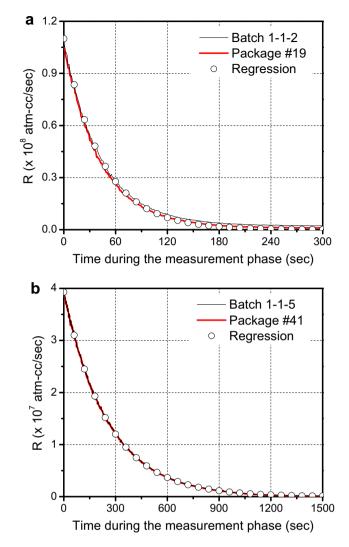


Fig. 8. Leak rate signals of a leaky package along with a batch that contains it: (a) Batch 1-1-2 and package #19 and (b) Batch 1-1-5 and package #41.

tribution of higher orders is negligible. Ignoring the higher order terms, Eq. (9) yields

$$\Omega \approx \Omega_1 + \Omega_2, \quad \Omega \lambda \approx \Omega_1 \lambda_1 + \Omega_2 \lambda_2, \quad \Omega \lambda^2 \approx \Omega_1 \lambda_1^2 + \Omega_2 \lambda_2^2, \tag{10}$$

The most ideal way to minimize the residual of regression is to match the first two lowest order terms, which gives

$$\Omega = \Omega_1 + \Omega_2$$
 and $\Omega \lambda = \Omega_1 \lambda_1 + \Omega_2 \lambda_2$ (11)

Substituting Eq. (11) into the third equation of (10) yields

$$\Omega_1 \Omega_2 (\lambda_1 - \lambda_2)^2 \approx 0 \tag{12}$$

The left-hand side of the above equation can be a representation of the *relative* residual of regression, a smaller value of which indicates higher possibility of false regression. Two possible cases to satisfy the equality are: (1) the true leak rates of two packages are identical, or (2) the product of two initial apparent leak rates is very small.

Fig. 9 illustrates the relative residual in Eq. (12) as a function of the true leak rate of the second package (L_2) when $L_1 = 10^{-7}$ atmcc/s (the corresponding λ_1 is 4.64 × 10⁻⁴). The cavity volume and test condition are identical to those used in the experiment. The corresponding initial leak rates were calculated from Eq. (2). The

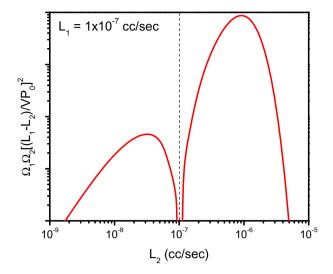


Fig. 9. Illustration of the residual of regression for two leaky packages (cavity volume = 2.156×10^{-4} cc, bombing time = 6 h, dwell time = 10 min, zero signal time = 150 s).

plot clearly shows two regimes where the regression is likely to fail; near $\lambda_1 \approx \lambda_2$ (or $L_1 \approx L_2$) and too small $\Omega_1 \Omega_2$ (when L_2 is too small or too high).

5.2. Simulation

A series of batch tests was simulated to illustrate the above limitations. As before, the true leak rate of the first package (L_1) was fixed as 10^{-7} atm-cc/s while that of the second package (L_2) was changed from $L_1/100$ to $26L_1$. The corresponding initial leak rates were obtained from Eq. (2). The result of regression analyses is summarized in Table 2. The prediction failed at the two regions identified earlier; when two true leak rates were identical and when one of leaky packages had a very low initial leak rate close to the measurement sensitivity. In the other cases, the regression predicted leaky packages very accurately (better than the experimental case due to the absence of measurement uncertainties).

The regression analysis was extended to three leaky packages. Three representative cases are presented in Table 3. In the first case the three packages with true leak rates in the median range (Table 2) are considered; the result is accurate in not only the number of leaky packages but also the true leak rates. The second case, which contained two packages with low initial leak rates, produced the accurate number of leaky packages but incorrect true leak rates information for two packages with low initial leak rates. In the third case a package with the true leak rate of 2.6×10^{-6} atm-cc/s, for which the regression failed to detect in the previous two-package simulation, replaced one of the packages. The regression failed again to capture this package.

This simulation study addresses clearly the application limit of the forward-stepwise regression analysis for the fine leak batch test; a feasible range of true leak rates is narrower than that of the individual test. In the simulations, the corresponding sensitivity limit of apparent leak rate was found to be approximately one

Table 3

Representative regression results for three leaky packages.

Accurate prediction			
Accurate prediction 10^{-7} (simulated)	1 (2 521)	0.2	F (4 201)
$L_2(\Omega_2) imes 10^{-7}$ (simulated)	1 (3.531)	0.2	5 (4.391)
		(0.807)	
$L_1(\Omega_1) imes 10^{-7}$ (predicted)	1.002	0.2004	5.009
	(3.531)	(3.531)	(4.391)
Difference (%)	0.2 (0.0)	0.2 (0.0)	0.2 (0.0)
Assumpts and istical of the number of leaf			
Accurate prediction of the number of leak		•	
$L_2\left(\Omega_2 ight) imes10^{-7}$ (simulated)	1 (3.531)	0.04	25
		(0.0651)	(0.0209)
$L_1(\Omega_1) \times 10^{-7}$ (predicted)	1.109	0.308	0.935
	(1.728)	(0.1531)	(1.722)
Difference (%)	10.9 (0.4)	. ,	96.3
()		(135.2)	(8139.2)
		(155.2)	(0155.2)
Failed prediction			
$L_2(\Omega_2) \times 10^{-7}$ (simulated)	1 (3.531)	0.04	26
2 (2)	((0.0651)	(0.0153)
$L_1(\Omega_1) \times 10^{-7}$ (predicted)	1.022	0.290	_
	(3.438)		
Difference (%)	• •	. ,	
Difference (%)	2.2 (2.6)	625.0	-
		(151.0)	

order higher than the measurement sensitivity, i.e., 10^{-9} atm-cc/s. It is speculated that this value will be a little bit larger for real tests due to measurement uncertainties.

6. Discussion: advanced formulations

When the true leak rates of two leaky packages are nearly identical, the regression analysis may recognize them as a single package with a doubled initial leak rate. It can be avoided by implementing "caps" for the initial leak rate as a function of the true leak rate, i.e., $c_1f(L_i) < \Omega_i < c_2f(L_i)$. The caps should be set based on a proper correlation such as Eq. (2). For example, if the failed case in Table 2 due to the identical true leak rate is solved using a constraint based on Eq. (2), the new prediction yields two leaky packages with true leak rates of 1.069×10^{-7} and 0.936×10^{-7} atm-cc/s, which are within an error range of 7%.

The existing models are not able to provide accurate relationship between true and initial leak rates for the whole range of measurable true leak rates [9]. It can be easily confirmed by applying Eq. (2) for the two leaky test packages; it produces the initial leak rates of 2.2×10^{-7} and 8.2×10^{-13} atm-cc/s, which are different from the true vales by 44% and five orders, respectively. This inaccuracy of the existing models warrants the development of an accurate physical model for gas transport through defective nano-channels.

The probability of incorrect prediction will increase as the number of leaky packages in a batch increases simply because there is higher probability for the batch to contain packages with nearly identical true leak rates or with very low initial leak rates. In order to reduce the risk of incorrect prediction a batch test can be designed in such a way that each batch should include only two or three leaky packages. In practice it may not be easy to assure the number of leaky packages in the batch. A multi-step test (or a hierarchical batch test) can be considered to cope with the problem [9],

Table 2

Regression results for two leaky packages ($L_1 = 1 \times 10^{-7}$ atm-cc/s, $\Omega_1 = 3.531 \times 10^{-7}$ atm-cc/s).

$L_2(\Omega_2) \times 10^{-7}$ (simulated) $L_1(\Omega_1) \times 10^{-7}$ (predicted) Difference (%)	0.01 (0.00475) 0.998 (3.534) 0.2 (0.1)	0.04 (0.0651) 1.002 (3.531) 0.2 (0.0)	0.2 (0.807) 1.002 (3.531) 0.2 (0.0)	1 (3.531) 1.002 (7.062) 0.2 (100.0)	5 (4.391) 1.002 (3.531) 0.2 (0.0)	25 (0.0209) 1.002 (3.531) 0.2 (0.0)	26 (0.0153) 1.003 (3.533) 0.3 (0.1)
$L_2(\Omega_2) \times 10^{-7}$ (predicted)	-	0.0401 (0.0651)	0.200 (0.807)	-	5.009 (4.391)	25.046 (0.0209)	-
Difference (%)	-	0.2 (0.0)	0.0 (0.0)	-	0.2 (0.0)	0.2 (0.0)	-

where a batch with more than three leaky packages is split into as many sub-batches as needed until the number of leaky packages in all sub-batches becomes less than 3.

7. Conclusions

An effective way of analyzing the fine leak rate batch test was proposed to infer the information of leaky metal-sealed MEMS packages within a batch. The forward-stepwise regression analysis scheme was employed to make the quantitative inference possible. The scheme was implemented for wafer-level hermetic MEMS packages. The true leak rates of the two packages were obtained from a single batch test data. The results were validated subsequently by extra batch tests as well as individual tests. Errors associated with the regression analysis were discussed to define the applicable domain of the proposed scheme. Advanced formulations were also suggested to extend the domain for a wider range of applications.

References

- Hablanian M. High-vacuum technology: a practical guide. New York: CRC Press; 1997. p. 493.
- [2] Knarr O. Industrial gaseous leak detection manual. New York: McGraw Hill; 1998. p. 97.
- [3] M. MIL-STD-883F. Test method standard microcircuits. US Department of Defense, 2004.
- [4] Goswami A, Han B, Ham S-J, Jeong B-G. Quantitative characterization of true leak rate of micro to nanoliter packages using a helium mass spectrometer. IEEE Trans Adv Packag 2009;32(2):440–7.
- [5] Howl DA, Mann CA. The back-pressurizing technique of leak testing. Vacuum 1965;15(7):347-52.
- [6] Goswami A, Han B. On ultra-fine leak detection of hermetic wafer level packages. IEEE Trans Adv Packag 2008;31(1):14–21.
- [7] Myers RH, Montgomery DC. Response surface methodology: process and product in optimization using designed experiments. New York: John Wiley & Sons; 1995.
- [8] Goswami A, Han B, Ham S-J, Jeong B-G. Quantitative characterization of true leak rate of micro to nanoliter packages using a helium mass spectrometer. IEEE Trans Adv Packag 2008.
- [9] Jang C, Youn BD, Ham S-J, Han B. Fine leak batch testing of multiple MEMS packages. In: Proc. 2009 SPIE, San Jose, CA, 720605, 2009.